

Theoretical Astrophysics and Cosmology

Spring Semester 2019 Prof. L. Mayer, Prof. J. Yoo

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Exercise Sheet 10

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Exercise 1: Linear evolution of the inflaton fluid quantities

a) Starting from the energy-momentum tensor of the inflaton field,

$$T_{\mu\nu} = \partial_{\mu}\phi \,\partial_{\nu}\phi - \frac{1}{2}g_{\mu\nu}\partial_{\rho}\phi \,\partial^{\rho}\phi - g_{\mu\nu}V(\phi) \,, \tag{1}$$

and consider scalar perturbations in conformal Newtonian gauge,

$$ds^{2} = -a^{2}(1+2\alpha)d\eta^{2} + a^{2}\delta_{\alpha\beta}(1+2\varphi)dx^{\alpha}dx^{\beta}, \qquad u^{\mu} = \frac{1}{a}(1-\alpha, -U^{,\alpha}), \qquad (2)$$

derive the linear-order fluid quantities:

$$\begin{aligned} \delta\rho &= \dot{\phi}\,\delta\dot{\phi} - \alpha\,\dot{\phi}^2 + \partial_{\phi}V\delta\phi \,,\\ \deltap &= \dot{\phi}\,\dot{\delta\phi} - \alpha\,\dot{\phi}^2 - \partial_{\phi}V\delta\phi \,. \end{aligned}$$
(3)

Hint: Recall that ρ and p are obtained from the stress-energy tensor as $\rho = T_{\mu\nu}u^{\mu}u^{\nu}$, and $p = \frac{1}{3}T_{\mu\nu}(g^{\mu\nu} + u^{\mu}u^{\nu})$. Be careful to not forget any linear order contributions! E.g., we have $\dot{\phi} = \partial_0 \phi$, but $\partial^0 \phi = g^{0\mu} \partial_\mu \phi \neq \dot{\phi}$.

b) Derive the equation of motion for the linear perturbation in the inflaton field:

$$\delta\ddot{\phi} + 3H\delta\dot{\phi} + \left(\partial_{\phi}^{2}V + \frac{k^{2}}{a^{2}}\right)\delta\phi = \dot{\phi}(\dot{\alpha} + \kappa) + (2\ddot{\phi} + 3H\dot{\phi})\alpha.$$

$$\tag{4}$$

Note that this equation is in Fourier space and k is the wavenumber. Also, recall that $\kappa = 3H\alpha - 3\dot{\varphi}$, and that $\alpha = -\varphi$ in a universe without anistropic stress.

Hint: You can start from the equation

$$0 = -(\partial_{\mu}\partial^{\mu}\phi + \Gamma^{\alpha}_{\alpha\mu}\partial^{\mu}\phi) + \frac{\partial V}{\partial\phi}, \qquad (5)$$

from Solution Sheet 7. Use the metric $ds^2 = -(1+2\alpha)dt^2 + a^2\delta_{\alpha\beta}(1+2\varphi)dx^{\alpha}dx^{\beta}$, with the cosmic time coordinate t.

Exercise 2: Linear evolution of the curvature perturbation

We define the gauge-invariant Mukhanov variable Φ , corresponding to the comoving gauge curvature perturbation, as

$$\Phi \equiv \varphi_v - \frac{K/a^2}{4\pi G(\rho + p)} \varphi_\chi \,, \tag{6}$$

where K is the extrinsic curvature and the gauge-invariant variables φ_v and φ_{χ} are defined in the lecture notes. Using the Einstein equations, derive the governing equations for the Mukhanov variable (eqs. (4.27)-(4.28) in the notes)

$$\Phi = \frac{H^2}{4\pi G(\rho+p)a} \left(\frac{a}{H}\varphi_{\chi}\right) + \frac{2H^2\Pi}{\rho+p},$$

$$\dot{\Phi} = -\frac{H}{4\pi G(\rho+p)} \frac{k^2 c_s^2}{a^2} \varphi_{\chi} - \frac{H}{\rho+p} \left(e - \frac{2k^2}{3a^2}\Pi\right),$$
(7)

where $e = \delta p - c_s^2 \delta \rho$, $c_s^2 = \dot{p}/\dot{\rho}$ is the speed of sound, and Π is the scalar part of the anisotropic stress tensor.

Furthermore, assuming a flat universe with K = 0, show that the second equation can be rearranged as (eq. (4.76) in the notes)

$$\dot{\Phi} = \Xi - \frac{H}{\rho + p} \frac{k^2}{a^2} \left(\frac{c_s^2}{4\pi G} \varphi_{\chi} - \frac{2}{3} \Pi \right), \qquad \Xi \equiv \frac{\dot{\rho} \,\delta p - \dot{p} \,\delta \rho}{3(\rho + p)^2} \,. \tag{8}$$

This means that the comoving gauge curvature perturbation is conserved in the super horizon limit $(k \to 0)$ if $\Xi = 0$ (i.e. $\delta p = \delta \rho \, dP/d\rho = \delta \rho \, \dot{P}/\dot{\rho}$, which sets the adiabatic initial condition).

Exercise 3: Inflation parameters for a given model

Work out the equations in Section 4.3.3 of the lecture notes for $\alpha = 2$.