

Issues in Standard Cosmology

$$\chi = \int_c^0 \frac{dt}{a} = \int_c^0 \frac{dln a}{aH} = \int_c^0 \frac{dz}{H_f} \quad -dt = a dr \quad dln a = -dln(1+z) = H dt$$

* particle horizon $\chi_h = \chi(t_0=0)$

for convergence of $a \rightarrow 0$. $aH \propto \sqrt{\Omega a^2 - K} \sim a^n \quad n < 0$

\therefore ok for mde, rde, but no Λ de

* $\rho \sim a^{-n}$, $a \sim t^{2/n}$
 $\chi \sim [t^{-2/n+1}]_{t_i}^{t_f}$

* $\rho \sim \Lambda$, $a \sim e^{Ht}$
 $\chi \sim [\frac{1}{H} e^{-Ht}]_{t_i}^{t_f}$

* event horizon $\chi_e = \chi(t_0 \rightarrow \infty)$

Yes for Λ de = closed, but no for flat or open

Horizon problem

mde $\chi = 6000^+ \text{ Mpc } (1+z)^{1/2} \rightarrow \chi_h(z=1100) \sim (80 \text{ h}^+ \text{ Mpc})$

comoving distance r to $z=1100 \rightarrow r = \int_0^z \frac{dz}{H} \sim 5820 \text{ h}^+ \text{ Mpc}$

$\rightarrow \theta \approx 1.8 \text{ degree} \approx 0.031$

no time for communication

Flatness problem

$\Omega_{tot} = \Omega_m + \Omega_r + \Omega_\Lambda + \dots$ without Ω_K .

$H^2 = \frac{8\pi G}{3} \rho - \frac{K}{a^2} \rightarrow -\frac{3K}{8\pi G \rho a^2} = \frac{3K}{8\pi G \rho} \dot{t} = \Omega_{tot}^+ \dot{t} \rightarrow \frac{\Omega_{tot}^+(t_{pr})}{\Omega_{tot}^+(t_0)} = \frac{(\rho a^2)_{pr}}{(\rho a^2)_{t_0}}$

today the deviation from unity $\Omega_{tot}^+ \approx 10$

but smaller by 10^{-60}

($K=0$ solved)

$\frac{\Omega_{tot}^+(t_{pr})}{\Omega_{tot}^+(t_0)} = \frac{(\rho a^2)_{pr}}{(\rho a^2)_{t_0}}$

$\sim \frac{T_0}{T_{pr}} \left(\frac{T_{pr}}{T_0} \right)^2 \sim 10^{-60}$

Monopole Problem (real problem?)

$T_{SUT} \sim 10^{16-15}$ GeV . topological defects from phase transition
 typically one defect per horizon of SUT $\rightarrow \rho_{today} \sim T_{SUT}^4 \rightarrow 10^{10}$

$$E_0 \sim 10^{-9}$$

↑

Structure formation problem

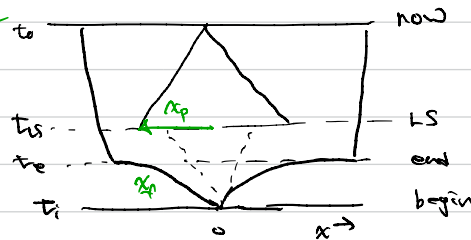
clusters of galaxies $M \sim 10^{15} M_\odot$. $R \sim 10$ Mpc comoving . $E_{grav} \sim 10^{-5}$ per unit proton mass
 should be in causal contact to collapse gravitationally $\rightarrow R < x_h \rightarrow$ only $z < 10^6$ (or BBN)
 but nothing in standard after BBN generates inhomogeneities of such high binding E .

\Rightarrow can be assumed as weird initial conditions

Inflation: accelerated expansion

$$x_p = \int_{t_i}^{t_0} \frac{dt}{a} \sim 3t_0 \quad \text{MDE}$$

$$x_p = \int_{t_i}^{t_f} \frac{dt}{a} = \frac{1}{a_0 H_{inf}} \left[e^{H_{inf} t} - 1 \right] \quad a = a_0 e^{H_{inf} (t-t_0)}$$



comoving horizon
 $\sim \frac{1}{aH}$ decrease
 \leftrightarrow larger in the past
 $\frac{d}{dt} \left(\frac{1}{aH} \right) < 0$

$$x_p \gg x_h \Rightarrow e^{H_{inf} t_0} > a_0 H_{inf} \cdot 3t_0 \sim \frac{T_e}{T_0} = 10^{26} \left(\frac{T_e}{10^{16} \text{ GeV}} \right) \Rightarrow H_{inf} \gtrsim 2 \ln 10 \approx 60$$

$$* \frac{\delta_i^2 - 1}{\delta_e^2 - 1} = \left(\frac{\delta_i}{\delta_e} \right)_e = \frac{a_e^2}{a_i^2} \gtrsim (10^{26})^2 = 10^{52}$$

flatness problem
 mitigated by expansion

$$\delta t \gtrsim 10^{-33} \text{ sec}$$

$$* \text{scale for monopole} \sim \frac{1}{a^2} \sim e^{-2H_{inf} t}$$

60 e-folding

Higgs EW transition ~ 200 GeV
 old inf. some GUT model
 with Higgs-like potential

Requirement

- * has to end. (Λ : cannot)
- * with large expansion. $T \rightarrow 0$ for ordinary radiation, if any.
 - \rightarrow reheating: needed to produce ordinary matter. inflation gone
 - hot enough for baryogenesis, but smooth enough for no monopole or defects

Models

* Old by Guth (original)

graceful exit problem

true vacuum: ordinary expansion

false vacuum: inflation \rightarrow large.

they have to collide to release the energy & stop \rightarrow too much inhomogeneities

* New by Linde, Steinhardt (similar but

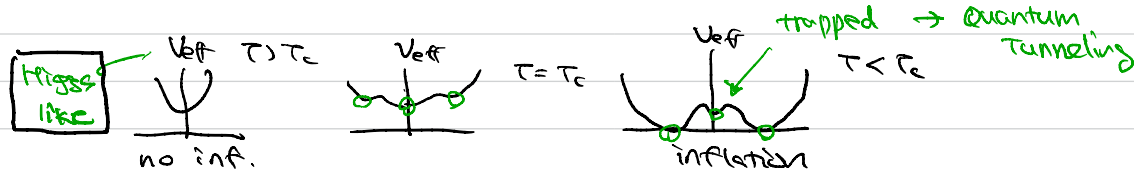
smooth 2nd-order phase transition, but slow-roll (\rightarrow minima $\gg M_{pl}$ unnatural).

also ϕ_{ini} should be $\ll M_{pl}$ otherwise stops too early \rightarrow fine-tuning in initial condition

* chaotic by Linde: $V = \frac{1}{2} m_p^2 \phi^2$
 or any power-law potential

theoretically some region satisfies $V \gg S$. slow-motion
 \rightarrow this region inflates \rightarrow our Univ.
 (other regions \times) but this requires $\phi_{ini} \gg M_{pl}$. Quantum grav?

* other notable models: R^2 Starobinsky, power-law infl. axion-like model



Higgs \neq inflaton
 \uparrow

