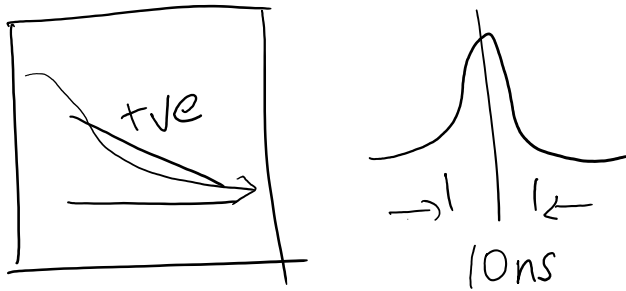


Noë Zimmermann winner!



★ Oral Exam Schedule is set via doodle

Recall: LAX Method:

$$\rho_i^{(n+1)} = \frac{1}{2}(\rho_{i+1}^{(n)} + \rho_{i-1}^{(n)}) - \frac{1}{2}C(\rho_{i+1}^{(n)} - \rho_{i-1}^{(n)})$$

$$C = \frac{\Delta t u}{\Delta x}$$

For linear advection equation, the flux is given by  $f(\rho) = \rho u$   
 Rewrite the above equation using fluxes:

$$\rho_i^{(n+1)} = \frac{1}{2}(\rho_{i+1}^{(n)} + \rho_{i-1}^{(n)}) - \frac{\Delta t}{2\Delta x}(f_{i+1}^{(n)} - f_{i-1}^{(n)})$$

Rewrite the 1<sup>st</sup> order Upwind Method:

$$\rho_i^{(n+1)} = \rho_i^{(n)} - \frac{\Delta t}{\Delta x}(f_i^{(n)} - f_{i-1}^{(n)}) \quad \text{for } u \geq 0$$

$$\rho_i^{(n+1)} = \rho_i^{(n)} - \frac{\Delta t}{\Delta x}(f_{i+1}^{(n)} - f_i^{(n)}) \quad \text{for } u < 0$$

Now for 1-D hydrodynamics, we have 3 conservation equations:

1. Conservation of mass:  $\bar{F} = \rho u$

2. Conservation of momentum:  $t = pu^x + P$   
 3. Conservation of energy:  $F = u(E + P)$

$P$  is the pressure.

Now let's use a "state vector"

$$\underline{U} = \begin{pmatrix} P \\ \rho u \\ E \end{pmatrix}$$

$$\underline{F}(\underline{U}) = [ \rho u, \rho u^2 + P, u(E + P) ]$$

little  $u$  is the velocity in the  $x$ -direction!

$$\frac{\partial \underline{U}}{\partial t} + \nabla \cdot \underline{F}(\underline{U}) = 0$$

Very general equation of "stuff".

We have 3 equations, but we have 4 unknowns:  $\rho, u, P, E$ !

$E =$  Kinetic Energy + Thermal Energy Density

$$= \frac{1}{2} \rho u^2 + \rho e \quad \begin{matrix} \text{(specific)} \\ \text{thermal energy density} \\ \text{specific thermal energy} \end{matrix}$$

Still have 4 unknowns...

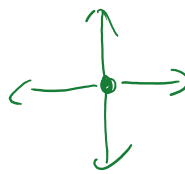
$\rho, u, P, e$

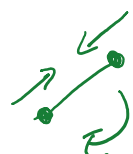
E.O.S. Equation of State Ideal Gas

$$e = \frac{P}{(\gamma - 1)\rho}$$

$$\gamma = \frac{f + 2}{f} \quad \text{where } f \text{ is the number of degrees of freedom}$$

↑ ... Monatomic Gas


 in 2-D Monoatomic Gas  
 $f = 2$   $\gamma = 2$


 in 1-D  $\gamma = 3$  we will consider this

$$E = \frac{1}{2} \rho u^2 + \frac{1}{2} P$$

$$\underline{F}(\underline{u}) = \left[ \rho u, \rho u^2 + P, \underbrace{u \left( \frac{1}{2} \rho u^2 + \frac{3}{2} P \right)}_{E+P} \right]$$

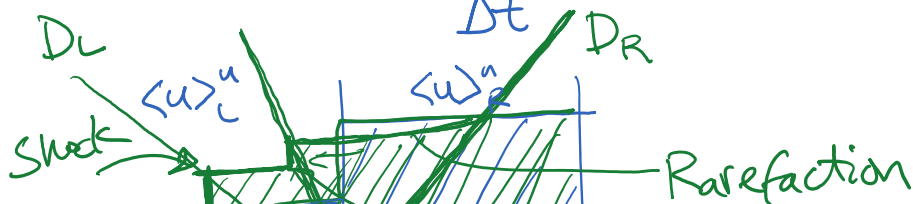
In finite volume let the values of  $\underline{u}$  in each cell be the average of the "material" in cell:

$$\langle u \rangle_i^n = \frac{1}{\Delta x} \int_{x_{i-1/2}}^{x_{i+1/2}} u(x, t^n) dx$$

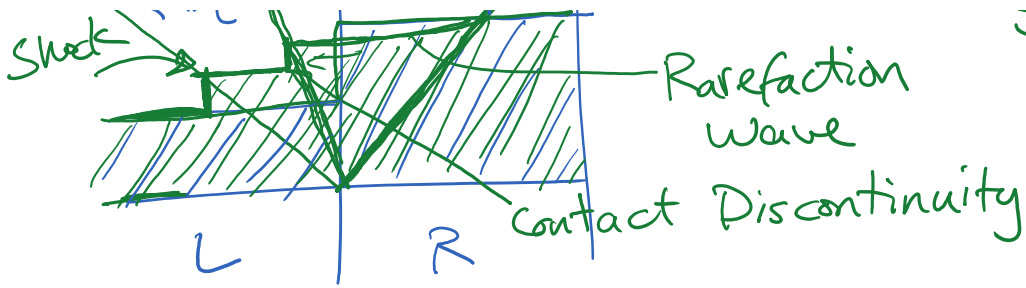
More complicated than for linear advection because information flows in more than one direction

$$\frac{\partial \underline{u}}{\partial t} + \nabla \cdot \underline{F}(\underline{u}) = 0$$

$$\frac{\langle u \rangle_i^{n+1} - \langle u \rangle_i^n}{\Delta t} + \frac{F_{i+1/2}^{n+1/2} - F_{i-1/2}^{n+1/2}}{\Delta x} = 0$$



3 Conserved Quantities



Conservation Quantities  
 $\Rightarrow$  3 waves  
 Characteristics  
 the 3 lines of information flow.

$$F(\underline{u}_L, \underline{u}_R) \leftarrow \text{Riemann Solver}$$

$D_L, D_R \leftarrow$  discontinuity speed of the left and right waves.

Approximate Riemann Solver:

$$D_L = -D_R \quad \text{simplification by assuming set to the maximum discontinuity speed, } D_{\max}$$

$$D = |u| + C_s \leftarrow \text{sound speed}$$

$$C_s = \sqrt{\gamma \frac{P}{\rho}} \quad \text{speed of sound}$$

$$D_{\max} = \max(D_L, D_R)$$

Lax-Friedrichs Riemann Solver

$$F_{i-\frac{1}{2}}^n = \frac{1}{2} [F(\underline{u}_i^n) + F(\underline{u}_{i-1}^n)] - \frac{1}{2} D_{\max} [\underline{u}_i^n - \underline{u}_{i-1}^n]$$

Back to our time-space centered approximation of the averages in cells:

Back to approximation of the averages in cells:

$$\underline{u}_i^{n+1} = \underline{u}_i^n - \frac{\Delta t}{\Delta x} \left[ F_{i+\frac{1}{2}}^{n+\frac{1}{2}} - F_{i-\frac{1}{2}}^{n+\frac{1}{2}} \right]$$

★ for  $n+\frac{1}{2}$  we could just use  $n$  for this  $\Rightarrow$  Method **[B]**

How could I get the flux at  $n+\frac{1}{2}$ :

We make a predictor step to  $n+\frac{1}{2}$  in order to calculate these fluxes:

$$* \underline{u}_i^{n+\frac{1}{2}} = \underline{u}_i^n - \frac{1}{2} \frac{\Delta t}{\Delta x} \left[ F_{i+\frac{1}{2}}^n - F_{i-\frac{1}{2}}^n \right] \text{ predictor step}$$

we follow this with an update step:

$$\underline{u}_i^{n+1} = \underline{u}_i^n - \frac{\Delta t}{\Delta x} \left[ F(* \underline{u}_i^{n+\frac{1}{2}}, * \underline{u}_{i+1}^{n+\frac{1}{2}}) - F(* \underline{u}_{i-1}^{n+\frac{1}{2}}, * \underline{u}_i^{n+\frac{1}{2}}) \right]$$

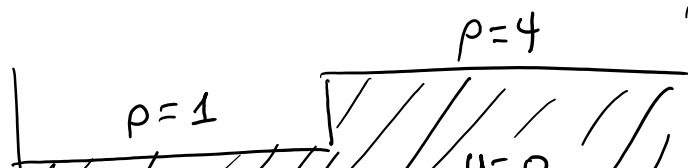
Method **[C]**

Method **[A]**: the original LAX Method

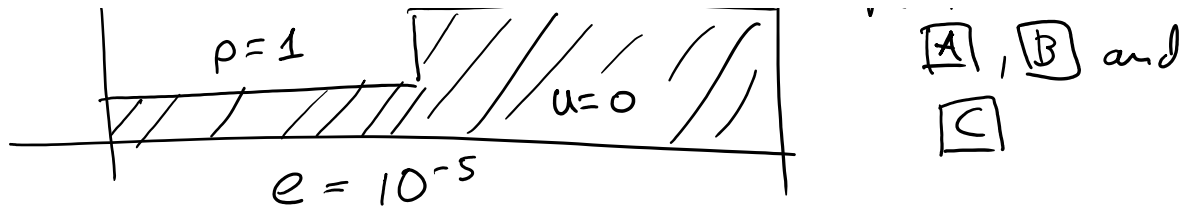
$$u_i^{n+1} = \frac{1}{2} (u_{i+1}^n + u_{i-1}^n) - \frac{\Delta t}{2\Delta x} (F_{i+1}^n - F_{i-1}^n)$$

\* No Riemann Solver involved!

Try is a "shock tube" problem

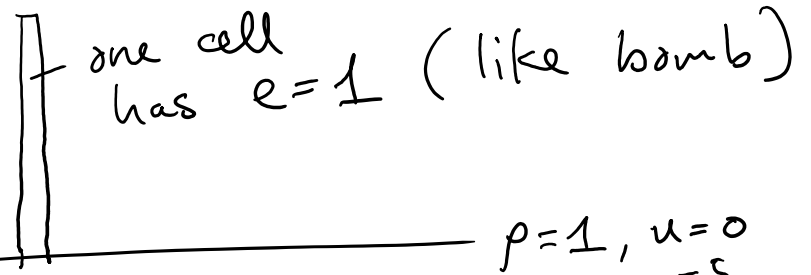


Methods **[A]**, **[B]** and **[C]**



→ solve for  $P$  given  $p$

## Sedov-Taylor Blast wave solution



$p=1, u=0$   
 $e=10^{-5}$

$\Delta t = ?$

Can use  $D_{max}$

like the Courant condition

$$D_{max} < \frac{\Delta x}{\Delta t}$$