

Finite Difference Method

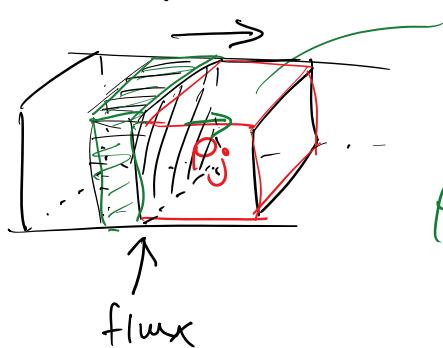
$$\frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} = 0$$

↑ ↑
operators on a grid

 "Jump"
 ↗ $\frac{\partial \rho}{\partial x}$ is not well defined

Conservation Laws lead to

Integral Equations



how much mass flows through the surfaces

$$\rho_j^{n+1} = \rho_j^n + \frac{\Delta t}{\Delta x} [f_{j-\frac{1}{2}} - f_{j+\frac{1}{2}}]$$

{ {

Integration over all fluxes
should be exactly 0.

Approximate These Fluxes

Linear Advection $f(\rho) = a \cdot \rho$ $a \geq 0$



$$\rho_j^{n+1} = \left[\left(\frac{a \cdot \Delta t}{\Delta x} \right) \rho_{j-1}^n + \left(1 - \frac{a \Delta t}{\Delta x} \right) \rho_j^n \right] c + \left[\left(1 - \frac{a \Delta t}{\Delta x} \right) \rho_{j-1}^n + \left(\frac{a \cdot \Delta t}{\Delta x} \right) \rho_j^n \right] (1-c)$$

Godunov's
Method

$$\rho_j^{n+1} = c \rho_{j-1}^n + (1-c) \rho_j^n \quad a \geq 0$$

$$\rho_j^{n+1} - \rho_j^n + c (\rho_j^n - \rho_{j-1}^n) = 0$$

$$\frac{U_j - U_{j-1}}{\Delta x}$$

1st order upwind method
Same in the case of linear advection.

Modified Equation

useless Method

$$\frac{\rho_j^{n+1} - \rho_j^n}{\Delta t} + a \frac{\rho_{j+1}^n - \rho_{j-1}^n}{2\Delta x} = 0$$

Recall: this is unstable

Taylor expand ρ_j^{n+1} to 2nd order:

(in time) $\rho_j^{n+1} = \rho_j^n + \Delta t \left(\frac{\partial \rho}{\partial t} \right) + \frac{\Delta t^2}{2} \left(\frac{\partial^2 \rho}{\partial t^2} \right) + \dots$

Taylor expand $\rho_{j+1}^n, \rho_{j-1}^n$ to 2nd order in x

$$\rho_{j+1}^n = \rho_j^n + \Delta x \left(\frac{\partial \rho}{\partial x} \right) + \frac{\Delta x^2}{2} \left(\frac{\partial^2 \rho}{\partial x^2} \right) + \dots$$

$$\rho_{j-1}^n = \rho_j^n - \Delta x \left(\frac{\partial \rho}{\partial x} \right) + \frac{\Delta x^2}{2} \left(\frac{\partial^2 \rho}{\partial x^2} \right) + \dots$$

$$\frac{\Delta t \left(\frac{\partial \rho}{\partial t} \right) + \frac{\Delta t^2}{2} \left(\frac{\partial^2 \rho}{\partial t^2} \right)}{\Delta t} + a \frac{(2\Delta x \left(\frac{\partial \rho}{\partial x} \right))}{2\Delta x} = 0$$

$$\Rightarrow \frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} = - \frac{\Delta t}{2} \left(\frac{\partial^2 \rho}{\partial t^2} \right)$$

$$\frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} = 0$$

$$\frac{\partial^2 \rho}{\partial t^2} + a \frac{\partial}{\partial x} \left(\frac{\partial \rho}{\partial t} \right) = 0$$

$$\frac{\partial^2 \rho}{\partial t^2} = - a^2 \frac{\Delta t}{2} \left(\frac{\partial^2 \rho}{\partial x^2} \right)$$

$$\frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} = - a^2 \frac{\Delta t}{2} \left(\frac{\partial^2 \rho}{\partial x^2} \right)$$

$$\frac{\partial \rho}{\partial t} + \left[a \frac{\partial \rho}{\partial x} \right] = -a^2 \frac{\Delta t}{2} \left(\frac{\partial^2 \rho}{\partial x^2} \right)$$

= 0 = D it is
 unstable
Diffusion

Ex. What is the modified equation for C.I.R. Scheme (1st order upwind)

2-D Advection

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \underline{u}) = 0$$

$$\underline{u} = \langle a, b \rangle \quad a, b > 0$$

$$\frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} + b \frac{\partial \rho}{\partial y} = 0$$

1st order upwind to solve this

$$\frac{P_{je}^{n+1} - P_{je}^n}{\Delta t} + a \frac{P_{je}^n - P_{je-1e}^n}{\Delta x} + b \frac{P_{je}^n - P_{je-1s}^n}{\Delta y} = 0$$

Stability Analysis shows that

$$C_a > 0 \quad C_a = \frac{a \Delta t}{\Delta x}$$

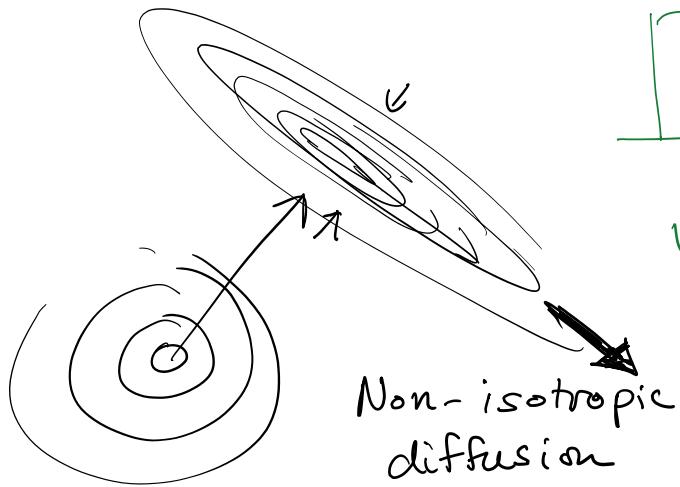
$$C_b > 0$$

$$\boxed{C_a + C_b \leq 1}$$

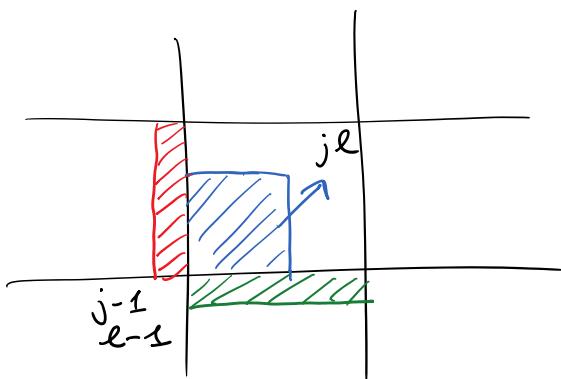
Sum of the 2 Courant numbers in x and y must be less than 1.

Modified Equation for this is interesting:

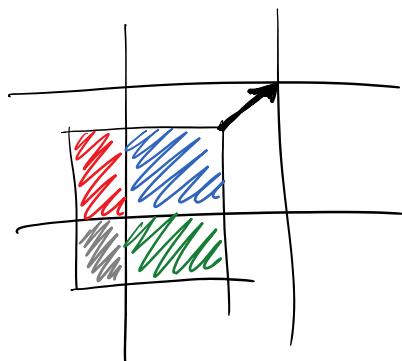
$$\frac{\partial \rho}{\partial t} + a \frac{\partial \rho}{\partial x} + b \frac{\partial \rho}{\partial y} = \frac{a \Delta x}{2} (1 - C_a) \frac{\partial^2 \rho}{\partial x^2} + b \Delta x r_1 \sim \frac{\partial^2 \rho}{\partial y^2}$$



$$\boxed{-ab\Delta t \frac{\partial^2 \rho}{\partial x \partial y}} \quad \begin{aligned} & \overline{z}(1-c_b) \overline{\frac{\partial \rho}{\partial yz}} \\ & \text{cross derivative term.} \\ & \text{won't preserve shape} \end{aligned}$$



Corner Transport upwind Method



$$\begin{aligned} \rho_{j,l}^{n+1} = & (1-c_a)(1-c_b)\rho_{j,l}^n \\ & + c_a(1-c_b)\rho_{j-1,l}^n \\ & + c_b(1-c_a)\rho_{j,l-1}^n \\ & + c_a c_b \rho_{j-1,l-1}^n \end{aligned}$$

Nice 2-step Method

$$\boxed{\begin{aligned} \rho_{j,l}^* &= (1-c_a)\rho_{j,l}^n + c_a \rho_{j-1,l}^n \\ \rho_{j,l}^{n+1} &= (1-c_b)\rho_{j,l}^* + c_b \rho_{j,l-1}^* \end{aligned}}$$

Stability is better as well

$$0 \leq c_a < 1, \quad 0 \leq c_b < 1$$

$$p(x, y) = A \exp \left[-\left(\frac{(x-x_0)^2}{2\sigma_x^2} + \frac{(y-y_0)^2}{2\sigma_y^2} \right) \right]$$

$$\sigma_x = \sigma_y$$

2-D C.I.R. vs CTU methods